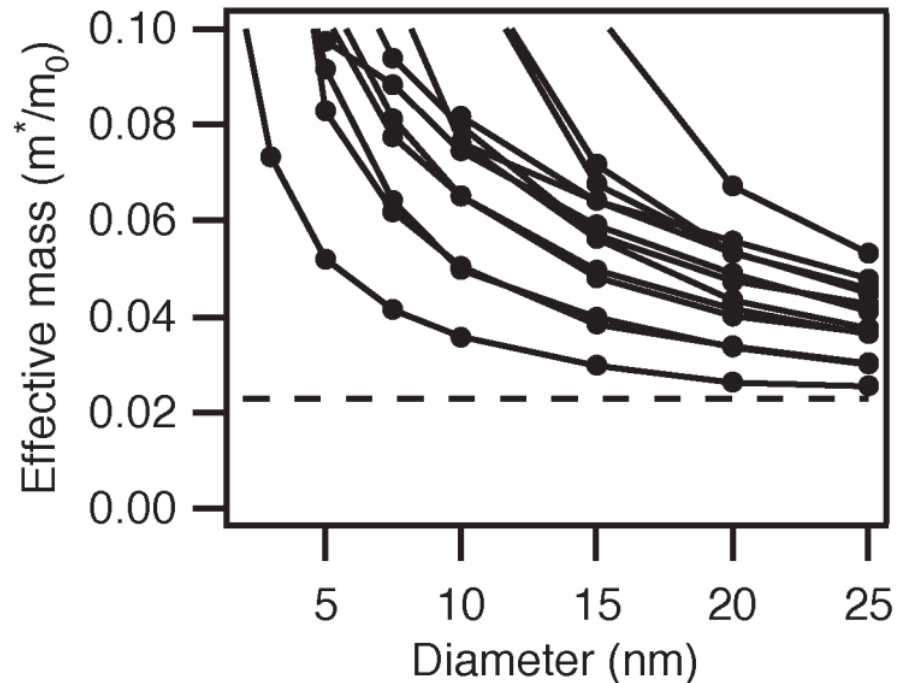
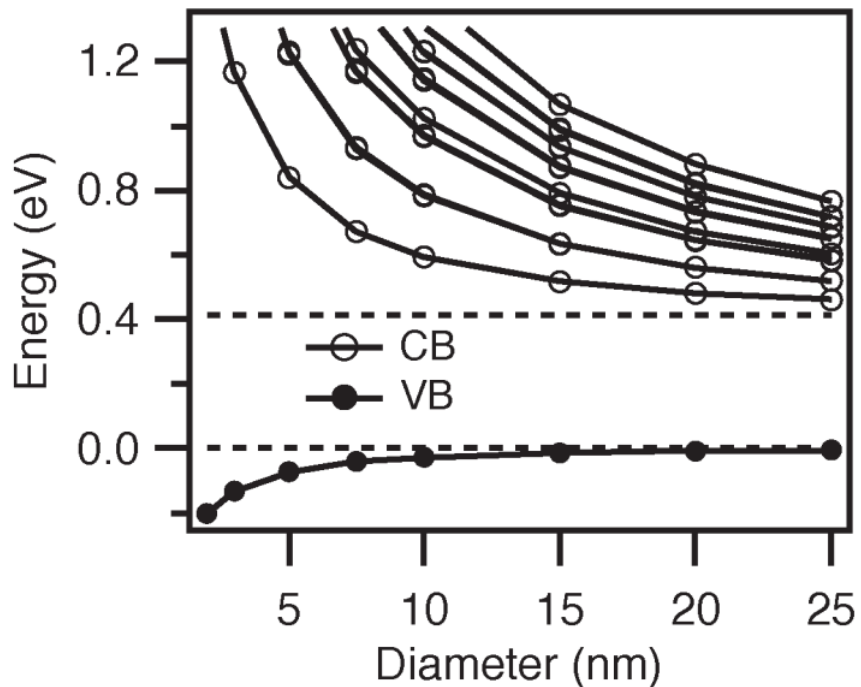


## Lecture 6 – 1D FETs

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- Nanowire MOSFETs
  - Nonparabolicity
- CNT Bandstructure
  - CNT FETs

# 1D Nanowires: Subbands spacing $\gg kT$



1D InAs nanowires – subband separation  $\ll kT$   
 $\sim 10$  nm diameter

Effective Mass:

$$E_{nm} = \frac{\hbar^2}{2m^*} \left( \frac{n^2}{W_z^2} + \frac{m^2}{W_y^2} \right)$$

# Nanowire MOSFET: Single Subband

$$n_L^+ = \frac{1}{L} \sum_{k>0} f_0 = \frac{N_{1D}}{2} F_{-\frac{1}{2}}(\eta_F)$$

$$n_L^- = \frac{1}{L} \sum_{k>0} f_0 = \frac{N_{1D}}{2} F_{-\frac{1}{2}}(\eta_F - U_D)$$

$$I_D^+ = \frac{1}{L} \sum_{k>0} v_x f_0 = \frac{2qkT}{h} F_0(\eta_F)$$

$$I_D^- = \frac{1}{L} \sum_{k>0} v_x f_0 = \frac{2qkT}{h} F_0(\eta_F - U_D)$$

$$v^+(0) = v_T \frac{F_0(\eta_F)}{F_{-\frac{1}{2}}(\eta_F)}$$

$$v^-(0) = v_T \frac{F_0(\eta_F - U_D)}{F_{-\frac{1}{2}}(\eta_F - U_D)}$$

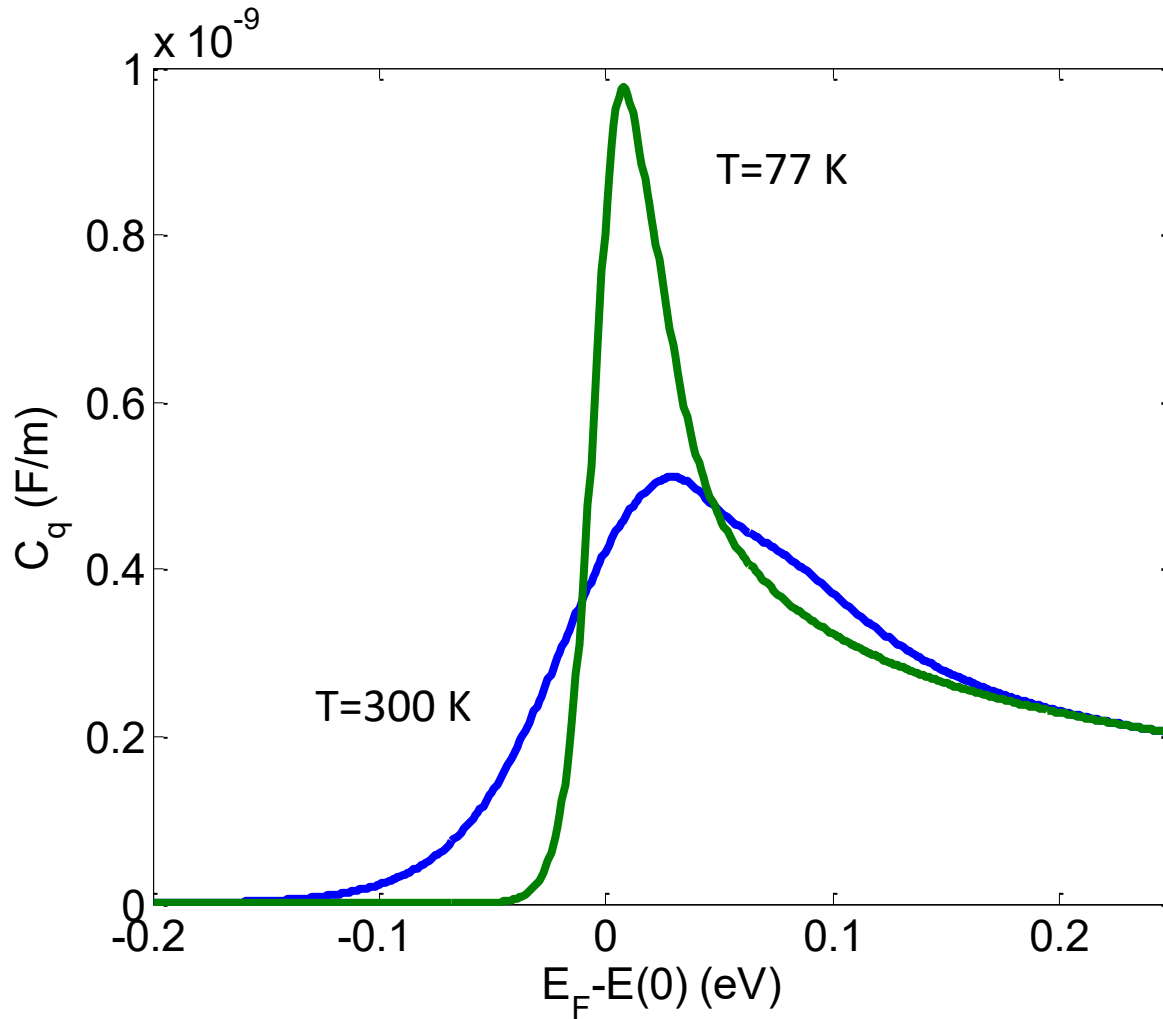
$$D_{1D} = \frac{\sqrt{2m^*}}{\pi\hbar} \frac{1}{\sqrt{E - E_C}}$$

$$N_{1D} = \sqrt{\frac{2m^*kT}{\pi\hbar^2}}$$

$$v_T = \sqrt{\frac{2kT}{\pi m^*}}$$

$$C_{ox} = \frac{2\pi\epsilon_r\epsilon_0}{\ln\left(\frac{t_{ox} + r_{wire}}{r_{wire}}\right)}$$

# Quantum Capacitance

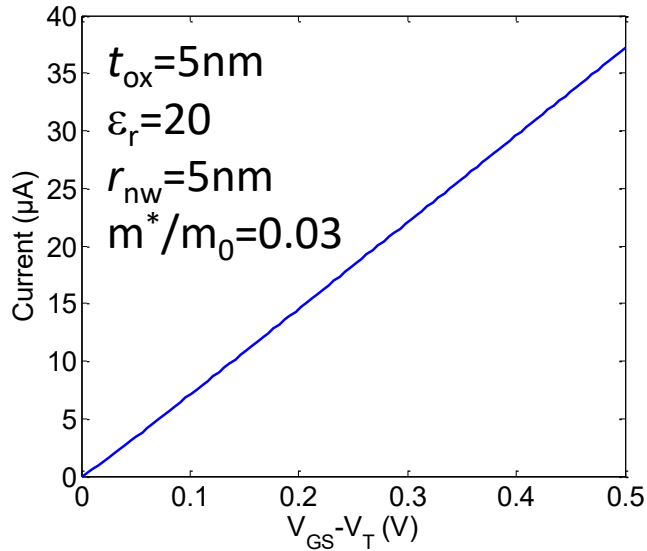


$$C_q = \frac{q^2 N_{1D}}{2kT} F_{-\frac{3}{2}}(\eta_F)$$

Degenerate:

$$C_q = \frac{\sqrt{2m^*} q^2}{\pi \hbar \sqrt{E_F - \varepsilon(0)}}$$

# Nanowire MOSFET – Degenerate conditions



A single subband InAs nanowire FET is operating very close to the ballistic limit!

Linear regime

$$I_{DS,lin} = \frac{2q^2}{h} V_{DS}$$

$$qV_{DS,sat} = \left[ -\frac{\sqrt{(2m^*)q^2}}{hC_{ox}} + \sqrt{\frac{2m^*q^4}{h^2C_{ox}^2} + q(V_{GS} - V_T)} \right]^2$$

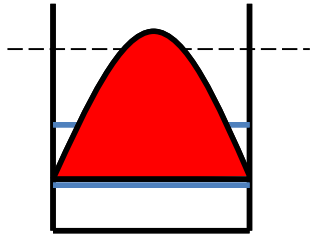
$$I_{D,sat} = \frac{2q^2}{h} V_{DS,sat}$$

$$C_q = \frac{\sqrt{(2m^*)q^2}}{\pi\hbar} \frac{1}{\sqrt{qV_{DS,sat}}}$$

$$I_{D,sat} = \frac{2q^2}{h} \frac{C_{ox}}{C_{ox} + C_q} (V_{GS} - V_T)$$

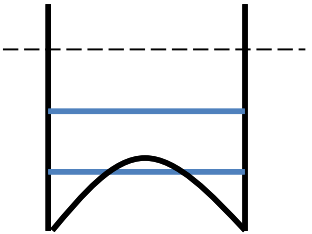
QCL Limit:  $I_{D,sat} = \frac{2q^2}{h} (V_{GS} - V_T)$

MOS Limit:  $I_{D,sat} = \frac{h}{4qm^*} C_{ox}^2 (V_{GS} - V_T)^2$

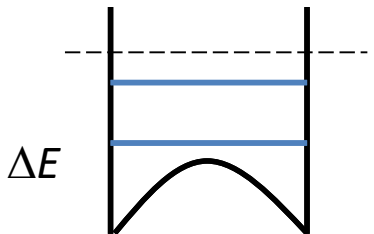
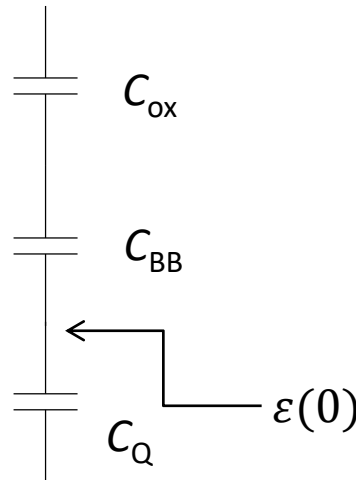


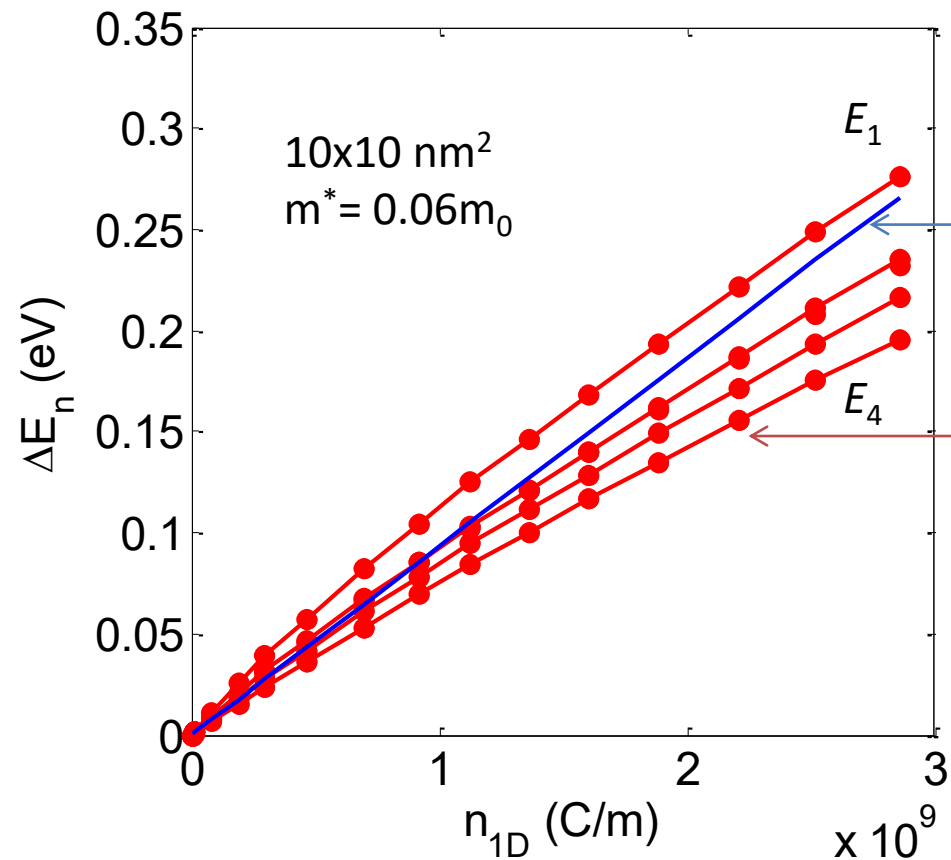
$$U(x, y) \approx \frac{16n_{1D}}{\pi^4} \frac{W_1 W_2}{W_1^2 + W_2^2} \sin\left(\frac{\pi x}{W_1}\right) \sin\left(\frac{\pi y}{W_2}\right)$$

$$\Delta E = \langle \psi_{1,1} | U(x, y) | \psi_{1,1} \rangle = \frac{\{0.164 - 0.29\}}{\epsilon_r \epsilon_0} \frac{64 W_1 W_2}{9 \pi^2 (W_1^2 + W_2^2)} n_{1D}$$



$$C_{BB} \approx 6.94 \epsilon_r \epsilon_0 \frac{(W_1^2 + W_2^2)}{W_1 W_2}$$



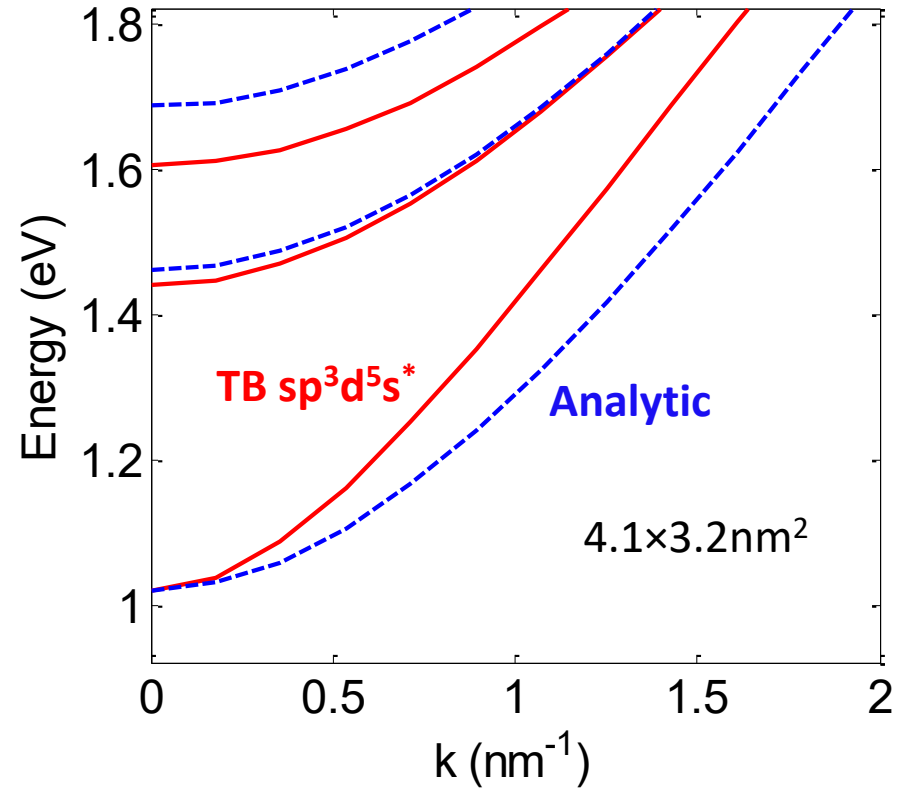
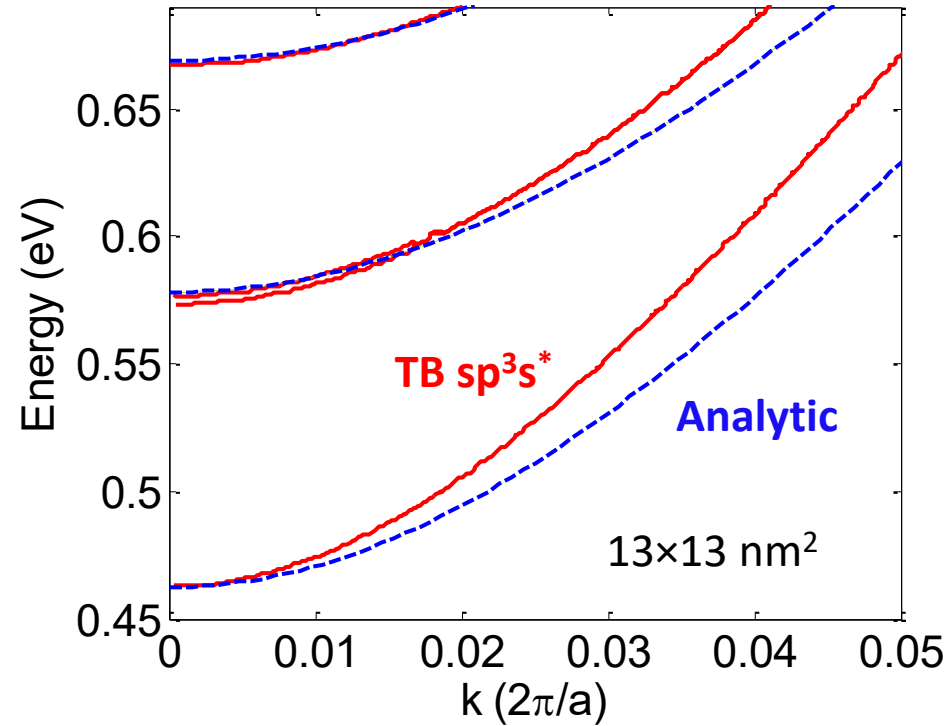


$$\Delta E = \langle \psi_{1,1} | U(x, y) | \psi_{1,1} \rangle$$

$$= \frac{0.2}{\epsilon_r \epsilon_0} \frac{64 W_1 W_2}{9 \pi^2 (W_1^2 + W_2^2)} n_{1D}$$

**Numerical Schrödinger-Possion**

# Nonparabolicity

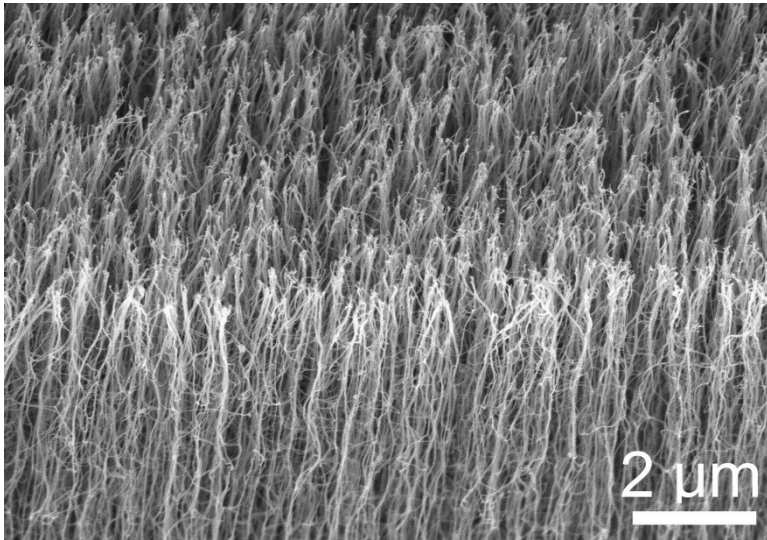
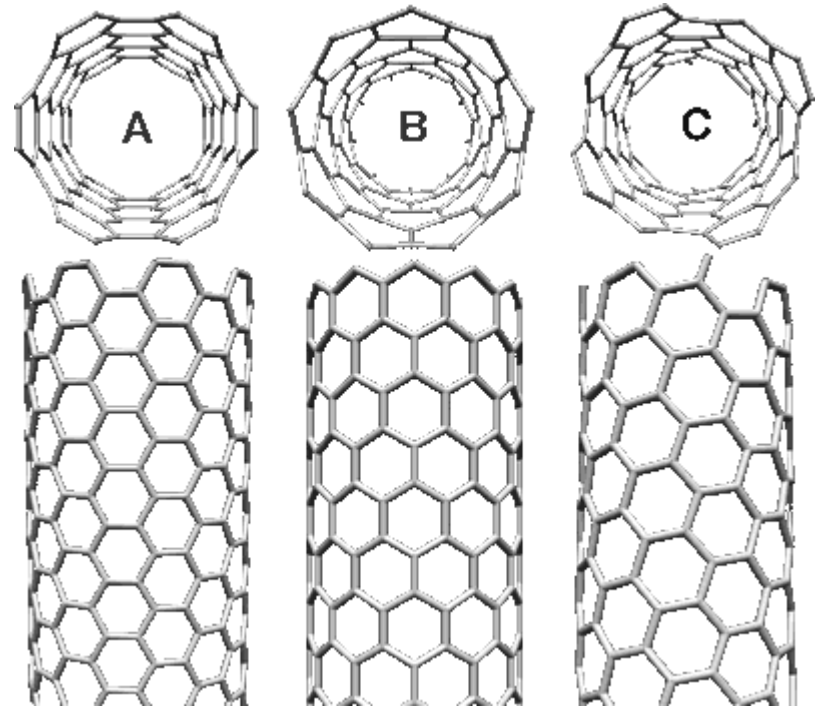
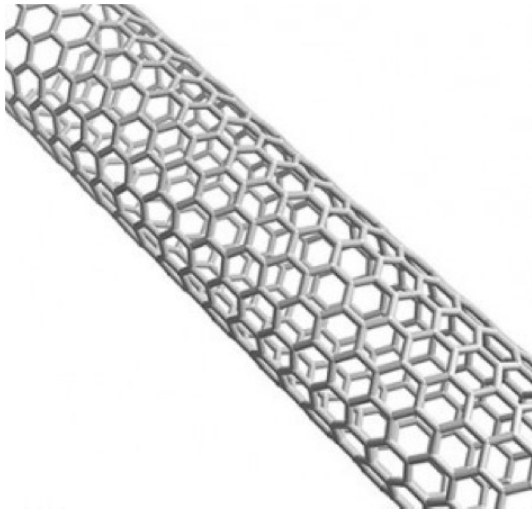


$$E(1 + \alpha_{n,m}E) = \frac{\hbar^2 k_x^2}{2m_{n,m}^*}$$

EMA:  $13 \times 13: \Delta E_{12} = 0.3 \text{ eV}$   
 $4 \times 3: \Delta E_{12} = 3 \text{ eV} (!!)$



# 1D Channels – Carbon Nanotubes



# Why CNTs?

**Traditional bulk semiconductors:**

Quantum Well/wire with surface roughness:

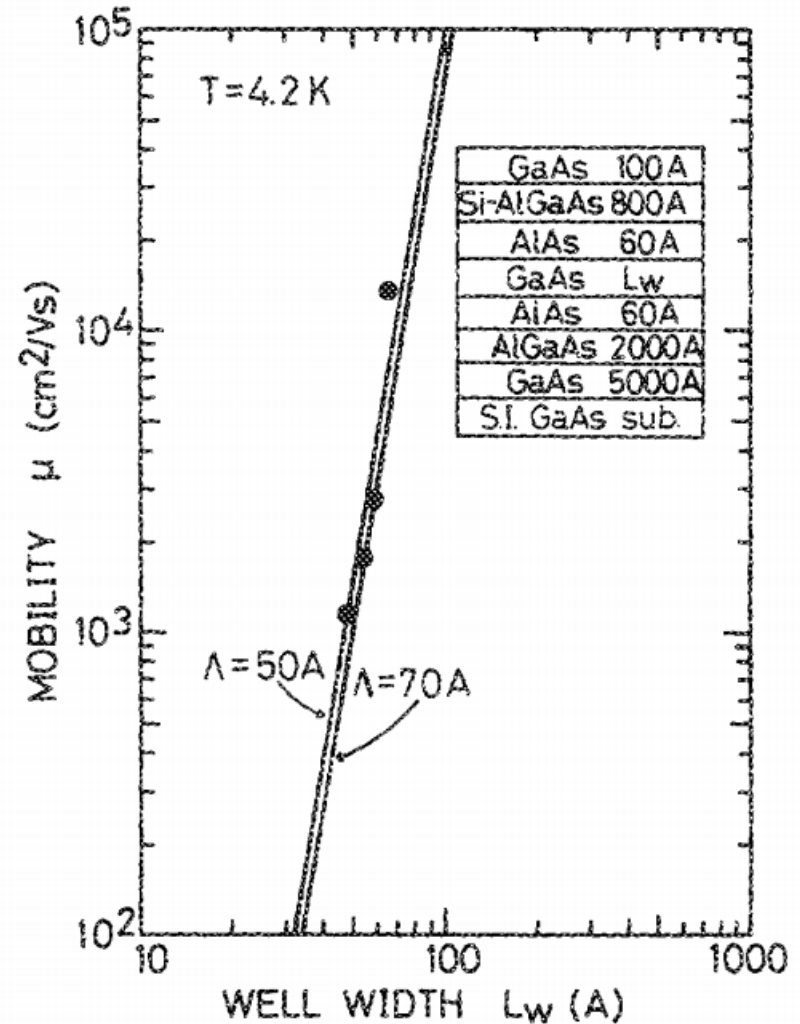
Mobility decreases as:  $\mu_{eff} \propto L_w^6$

Strong asymmetry between electron/hole effective masses

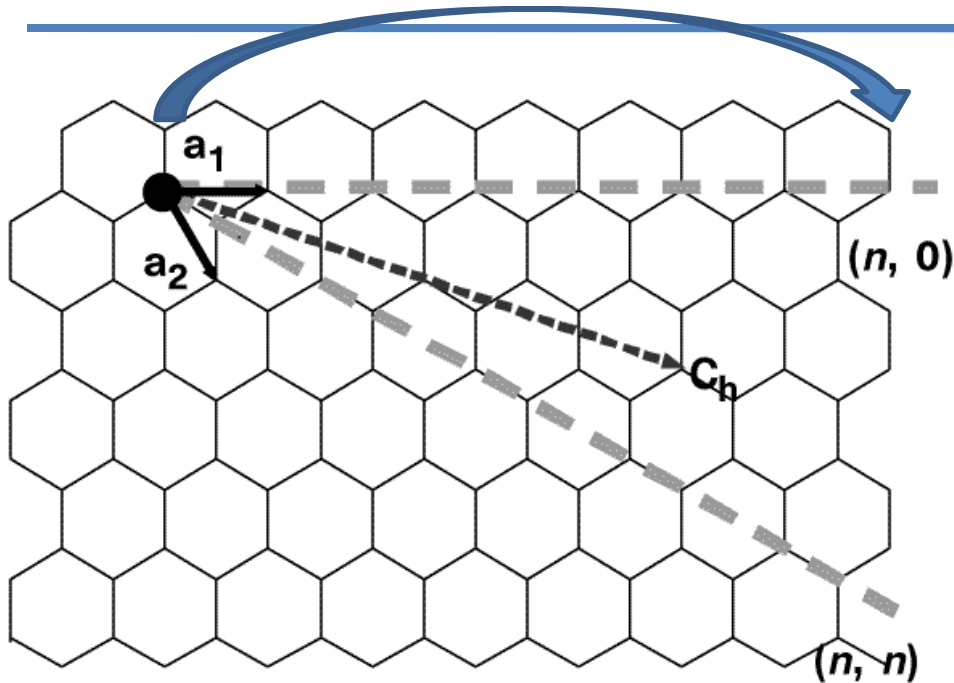
**CNTs (& graphene, 2D materials):**

No 'dangling' bonds or surface reconstruction. Naturally atomically flat.

**H. Sakaki, T. Noda, K. Hirakawa, M. Tanaka and T. Matsusue, Applied Physics Letters 51 (23), 1934-1936 (1987).**



# CNT formation & bandstructure



$$\hat{C} = \hat{a}_1 n + \hat{a}_2 m$$

$$d = \frac{\sqrt{3}a_{cc}}{\pi} \sqrt{n^2 + m^2 + nm}$$

**Metallic** if:  $(n-m)$  is a multiple of 3  
**Semiconductor** otherwise

Armchair:  $(n,n)$  always metallic

Zig-zag:  $(n,-n)=(n,0)$

Chiral:  $(n,m)$

A CNT has a valley degeneracy = 2

$$E_G = \frac{2a_{cc}t}{d} \approx \frac{0.8}{d(\text{in nm})} \text{ eV}$$

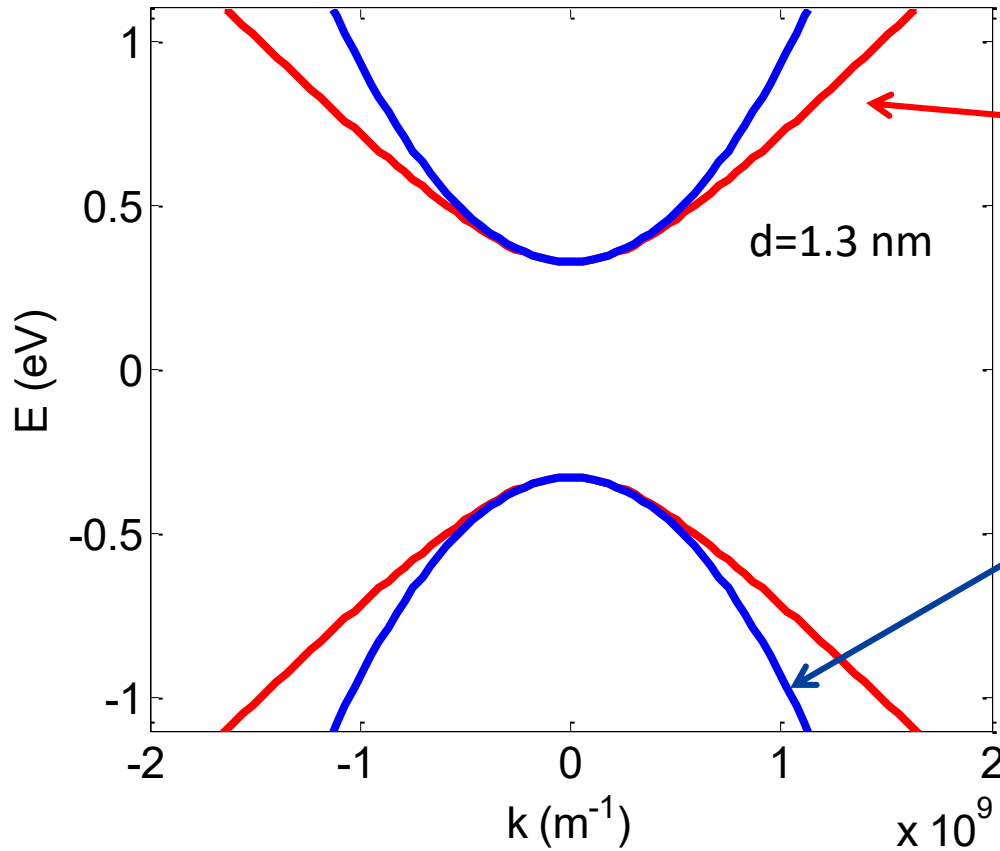
$$a_{cc} = 1.42 \text{ \AA}$$

$$t = 3.0 \text{ eV}$$

$$E(k) \approx \pm \frac{E_G}{2} \sqrt{1 + \left(\frac{3kd}{2}\right)^2}$$

$$E_{Gi} = \frac{2a_{cc}t}{d} \times \frac{6i - 3 - (-1)^i}{4}$$

# CNT bandstructure



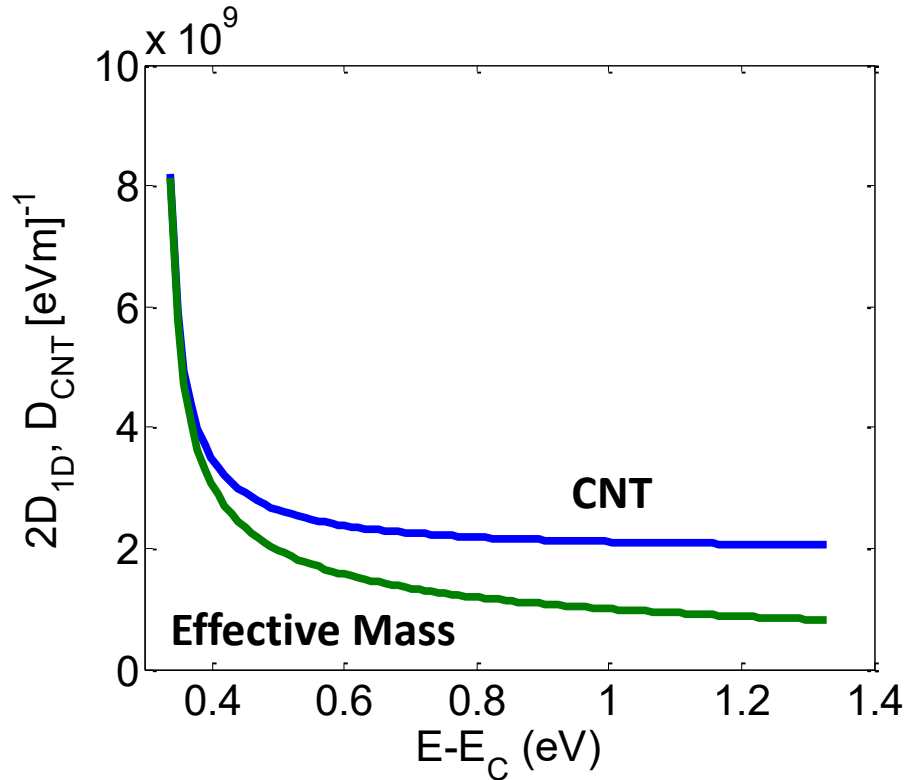
$$E(k) \approx \pm \frac{E_G}{2} \sqrt{1 + \left(\frac{3kd}{2}\right)^2}$$

$$\frac{m^*}{m_0} = \frac{0.08}{d(\text{nm})}$$

$$E(k) \approx \pm \left( \frac{E_G}{2} + \frac{\hbar^2 k^2}{2m^*} \right)$$

Non-parabolicity is strong for CNTs. Effective mass only valid around  $\sim 0.1$  eV.

# Effect of nonparabolicity



CNT DOS

$$D(E) = \frac{D_0 |E|}{\sqrt{E^2 - \frac{E_G^2}{4}}}, |E| \geq \frac{E_G}{2}$$

Effective Mass

$$D_{1D} = \frac{\sqrt{2m^*}}{\pi\hbar} \frac{1}{\sqrt{E - E_C}}$$

Quantum capacitance:  $C_q$  – DOS at  $E_F$

Ignoring band structure effects will underestimate  $C_q$ !